

**Abstract for the lecture on March 8 2017**

Contents from Torsten Carleman's book *Sur les equations integrales singulier a noyaux réell et symétrique* will be exposed. The major results deal with spectral resolutions of unbounded self-adjoint operators on complex Hilbert spaces. I will describe some applications to specific problems such as PDE-equations and Stieltjes' moment problem.

### An example

Following page 174-185 in Carleman's cited book we announce a result about *propagation of sound* which amounts to find a function  $u = u(x, y, z, t)$  where  $t$  is a positive time variable and  $(x, y, z)$  are points in  $\mathbf{R}^3$  which stay outside a bounded domain  $\Omega$  whose boundary  $S$  is of class  $C^1$ . Put  $U = \mathbf{R}^3 \setminus \Omega$  and consider the family  $\mathcal{F}$  of  $C^2$ - functions  $f(x, y, z)$  in  $U$  such that  $f$  and the Laplacian  $\Delta(f)$  both belong to  $L^2(U)$  and the normal derivative

$$\frac{\partial f}{\partial \mathbf{n}} = 0$$

along  $S$ . Now one seeks a function  $u(x, y, z, t)$  defined in  $U \times \mathbf{R}^+$  satisfying the wave equation

$$\frac{\partial^2 u}{\partial t^2} = \Delta(u)$$

and the following three conditions

$$(1) \quad \frac{\partial u}{\partial \mathbf{n}}(p) = 0 \quad p = (x, y, z) \in S \text{ \& } t > 0$$

$$(2) \quad u(p, 0) = f_0(p) \quad p \in U$$

$$(3) \quad \frac{\partial u}{\partial t}(p, 0) = f_1(p) \quad p \in U$$

Apart from existence and uniqueness of a solution to this boundary value problem for every pair  $f_0, f_1$  one may also ask if the motion tends to zero as  $t \rightarrow +\infty$  inside every bounded part of  $U$ . In other words if

$$(*) \quad \lim_{t \rightarrow +\infty} \left| \frac{\partial u}{\partial x}(p, t) \right| + \left| \frac{\partial u}{\partial y}(p, t) \right| + \left| \frac{\partial u}{\partial z}(p, t) \right| = 0$$

when  $p$  stays in a bounded subset of  $U$ .

The proof of these results rely upon a study of the Laplace operator  $\Delta$  acting on the space  $\mathcal{F}$ . Here  $\Delta$  is only densely defined on the Hilbert space  $L^2(U)$  and in addition unbounded. So one must employ the general theory created in Carleman's work from 1923. A crucial step towards the existence of a solution to the boundary value problem is that when  $f \in \mathcal{F}$  then the Dirichlet integral

$$\iint_U (\partial^2 f / \partial x^2 + \partial^2 f / \partial y^2 + \partial^2 f / \partial z^2) dx dy dz < +\infty$$

Using this Carleman proved that  $-\Delta$  is a densely defined self-adjoint operator on  $\mathcal{F}$  whose spectrum is confined to  $\mathbf{R}^+$ . The requested solution to the boundary value problem for a given pair  $f_0, f_1$  is given by

$$u(p) = \int_0^\infty \cos \sqrt{\lambda} \cdot t \, d\lambda \int_U \Theta(p, q, \lambda)(f_0)(q) \, dq + \int_0^\infty \frac{\sin \sqrt{\lambda} \cdot t}{\sqrt{\lambda}} \int_U \Theta(p, q, \lambda)(f_1)(q) \, dq$$

where  $\Theta(p, q, \lambda)$  is the spectral function attached to the self-adjoint operator  $-\Delta$ .

Finally the proof of (\*) employs spherical functions in  $\mathbf{R}^3$  which reduces the proof of the absolute continuity of the spectral function  $\Theta$  to a result concerned with a certain second order ODE on the real line which is presented on page 184-185 in Carleman's cited book.

## Introduction.

### Hyperbolic PDE-equations.

**Introduction.** We shall study boundary value problems for linear symmetric hyperbolic systems. The proof of our main result in Theorem 0.1 teaches the usefulness of regarding densely defined unbounded linear operators on Hilbert spaces. The proof of Theorem 0.1 involves quite technical steps. For the reader's convenience we therefore include a self-contained proof in the restricted case for systems in two variables and a single scalar function. Here the Hilbert space methods are transparent, and at the same time the methods which are used in this restricted situation are crucial, i.e. the general case is verbatim the same except for various technical steps. Let us first recall some background about hyperbolic equations.

A classic result due to Hadamard gives a vanishing principle for well-posed boundary value problems. With coordinates  $(x, s) = (x_1, \dots, x_n, s)$  in  $\mathbf{R}^{n+1}$  we consider a differential operator of the form

$$Q(x, s, \partial_x, \partial_s) = \partial_s^p + \sum_{\nu=0}^{p-1} P_\nu(x, s, \partial_x) \cdot \partial_s^\nu$$

where  $\{P_\nu(x, s, \partial_x)\}$  are differential operators which are independent of  $\partial_s$  and coefficients in  $C^\infty(\mathbf{R}^{n+1})$  which in general are complex-valued. The Cauchy problem is well posed in Hadamard's sense if there to every  $f(x) \in C^\infty(\mathbf{R}^n)$  exists a unique  $C^\infty$ -function  $g(x, s)$  in the half-space  $\{s \geq 0\}$  such that  $Q(g)(x, s) = 0$  when  $s > 0$  and on  $s = 0$  one has

$$\partial_s^\nu(g)(x, 0) = 0 : 0 \leq \nu \leq p-1 \quad : \quad \partial_s^p(g)(x, 0) = f(x)$$

Under the hypothesis that Cauchy's problem is well-posed one has:

**Hadamard's Theorem.** *If  $K$  is a compact subset in  $\mathbf{R}^{n+1}$  there exists a compact set  $K$  in the  $x$ -space such that there unique solution  $f$  whose Cauchy data on  $s = 0$  vanishes on  $k$  must vanish on  $K$ .*

This result follows easily from Baire's category theorem. The reader may consult [D-S: Volume 2: page 1649-1652] for details.

Conditions for a PDE-operator to be hyperbolic is an extensive subject. The reader may consult the text-book by Petrowsky for examples of hyperbolic, elliptic and parabolic equations.

**An ill-posed equation.** Let  $n = 1$  and consider the  $2 \times 2$ -matrices

$$A_1(x, s) = \begin{pmatrix} -e^{-x} & 0 \\ 0 & 1 \end{pmatrix} \quad : \quad B(x, s) = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \quad :$$

One seeks pairs of functions  $(f_1(x, s), f_2(x, s))$  which satisfy the first order system:

$$\frac{\partial f_1}{\partial s} = -e^{-x} \cdot \frac{\partial f_1}{\partial x} + f_2 \quad : \quad \frac{\partial f_2}{\partial s} = \frac{\partial f_2}{\partial x}$$

For any function  $C^\infty$ -function  $h$  of a single variable we see that

$$f_1(x, s) = h(e^x - s) \quad : \quad f_2 = 0$$

solves the system above and here  $f_1(x, 0) = h(x)$ . Consider the singleton set  $\{0, 1\}$  in  $\mathbf{R}^2$ . Then  $f(A) = 1$  for all  $h$ -functions such that  $h(0) = 0$ . Fix a test-function  $\phi(t)$  on the real  $t$ -line where  $\phi(0) = 1$  while  $\phi(t) = 0$  if  $|t| \geq 1/2$ . For every positive integer  $N$  we take  $h(t) = \phi(Nt)$ . If the Cauchy problem is well posed we get the unique solution with  $f_1 = h(s - e^x)$  and now

$$f_1(x, 0) = \phi(e^{Nx})$$

Since the support of  $\phi$  is contained in  $[-1/2, 1/2]$  we see that  $f_1(x, 0) \neq 0$  entails that

$$e^{Nx} \leq 1/2 \implies x \leq -N \cdot \log 2$$

Since  $N$  can be arbitrary large this violates Hadamard's vanishing principle and hence the Cauchy problem for this system is not well-posed.

### The wave operator.

With two real variables we consider the PDE-operator

$$P = \frac{\partial^2}{\partial x_1^2} - \frac{\partial^2}{\partial x_2^2}$$

To each pair of functions  $g(x_1), h(x_1)$  one seeks  $f(x_1, x_2)$  such that  $P(f) = 0$  and Cauchy's boundary value conditions:

$$(i) \quad f(x_1, 0) = g(x_1) \quad : \quad \frac{\partial f}{\partial x_2}(x_1, 0) = h(x_1)$$

This boundary value problem corresponds to a first order system where one seeks a pair  $f_1(x_1, x_2)$  and  $f_2(x_1, x_2)$  such that

$$(ii) \quad \frac{\partial f_1}{\partial x_1} = -\frac{\partial f_1}{\partial x_2} + f_2 \quad : \quad \frac{\partial f_2}{\partial x_1} = \frac{\partial f_2}{\partial x_2}$$

with boundary values

$$(iii) \quad f_1(x_1, 0) = g(x_1) \quad : \quad f_2(x_1, 0) = g'(x_1) + h(x_1)$$

**Exercise.** Show that if  $f$  solves (i) then the boundary value system is solved by the pair

$$f_1 = f \quad : \quad f_2 = \frac{\partial f}{\partial x_1} + \frac{\partial f}{\partial x_2}$$

Show also that if  $f_1, f_2$  solves the system then  $f = f_1$  solves the original equation. Next, consider the matrices

$$A = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} \quad : \quad B = \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

Then the system can be written in matrix form as

$$\frac{\partial}{\partial x_1} \begin{pmatrix} f_1 \\ f_2 \end{pmatrix} = A \frac{\partial}{\partial x_2} \begin{pmatrix} f_1 \\ f_2 \end{pmatrix} + B \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}$$

This clarifies that the original equation is a first order symmetric system to be described in § xx,

**Higher order systems.** Consider systems with  $n+1$  many variables  $x_1, \dots, x_n, s$  where the variable  $s$  is distinguished. In a first order scalar system one seeks a function  $f(x, s)$  such that

$$\frac{\partial f}{\partial s} = \sum_{j=1}^{j=n} a_j(x, s) \cdot \frac{\partial f}{\partial x_j} + b(x, s)$$

satisfying the boundary condition

$$f(x, 0) = g(x)$$

where the  $g$  function is given in the  $n$ -dimensional  $x$ -space. In a vector-valued system of order  $m \geq 2$  the  $a$ -functions are replaced by  $m \times m$ -matrices and  $b$  by some  $m \times m$ -matrix. Here one seeks a vector-valued function  $f = (f_1, \dots, f_m)$  such that

$$(*) \quad \frac{\partial}{\partial s} \begin{pmatrix} f_1 \\ \dots \\ f_m \end{pmatrix} = \sum_{j=1}^{j=n} A_j(x, s) \cdot \frac{\partial}{\partial x_j} \begin{pmatrix} f_1 \\ \dots \\ f_m \end{pmatrix} + B(x, s) \begin{pmatrix} f_1 \\ \dots \\ f_m \end{pmatrix}$$

The boundary conditions are expressed by an  $m$ -tuple of functions  $\{g_\nu(x)\}$  such that  $f_\nu(x, 0) = g_\nu(x)$  hold for each  $\nu$ . The  $A$ -matrices and the  $B$ -matrix are in general complex valued. A famous example is:

**Maxwell's equations of electrodynamics.** Let  $n = m = 3$  where  $\{A_\nu\}$  are constant  $3 \times 3$ -matrices

$$A_1 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix} : A_2 = \begin{pmatrix} 0 & 0 & i \\ 0 & 0 & 0 \\ -i & 0 & 0 \end{pmatrix} : A_3 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Moreover  $B = 0$  and one seeks a vector-valued function  $f_1, f_2, f_3$  which satisfies (\*) and boundary value conditions  $f_\nu(x, 0) = g_\nu(x)$

**The symmetric case.**

The system (\*) is symmetric if the matrices  $\{A_\nu(x, s)\}$  are Hermitian for each  $1 \leq \nu \leq n$ . No special condition is imposed on  $B$ , i.e. it can be an arbitrary complex  $m \times m$ -matrix. Notice that the  $A$ -matrices in Maxwell's equations are hermitian. The following result is due to Friedrichs:

**0.1 Theorem.** *Assume that the  $A$ -matrices are hermitian and that the matrix elements of  $A_1, \dots, A_n$  and  $B$  are bounded functions in  $\mathbf{R}^{n+1}$ . Then Cauchy's boundary value problem has a unique  $C^\infty$  solution  $f = (f_1, \dots, f_m)$  for every  $m$  tuple  $g = (g_1, \dots, g_m)$  of  $C^\infty$ -functions in the  $n$ -dimensional  $x$ -space.*

**Remark.** The example after Hadamard's result above shows that this boundedness is needed in order that the Cauchy problem is well-posed.

**0.2 On the uniqueness.** Let us illustrate why the condition that the  $A$ -matrices are hermitian gives a certain vanishing principle. Let  $\Omega$  be a bounded open set in the  $n$ -dimensional  $x$ -space and  $-s_* < s < s_*$  is some open  $s$ -interval. Let  $\{A_j(x, s)\}$  be Hermitian  $m \times m$ -matrices whose elements as well as their first order partial derivatives are bounded  $C^\infty$ -functions in  $\Omega \times (-s_*, s_*)$ . Similarly, assume that  $B(x, s)$  is an  $m \times m$ -matrix whose elements are bounded  $C^\infty$ -functions in  $\Omega$ . Let  $f = (f_1, \dots, f_m)$  be a vector-valued solution to the system (\*) where

$$(0.2.1) \quad f(x, 0) = 0 \quad : x \in \Omega$$

and assume there is a compact subset  $K$  of  $\Omega$  and  $f = 0$  in  $(\Omega \setminus K) \times (-s_*, s_*)$ . Then (0.2.1) entails that  $f = 0$  in  $\Omega \times (-s_*, s_*)$ . To prove this we introduce the function

$$J(s) = \int_{\Omega} |f(x, s)|^2 dx$$

where  $|f|^2 = \sum f_\nu \cdot \bar{f}_\nu$ . Taking the derivative with respect to  $s$  we get

$$(i) \quad \frac{dJ}{ds} = 2 \cdot \Re \sum_{j=1}^{j_m} \int \partial_s(f_\nu) \cdot \bar{f}_\nu dx$$

Since  $f$  satisfies (\*) we have

$$\sum_{j=1}^{j_m} \int \partial_s(f_\nu) \cdot \bar{f}_\nu = \sum_{j=1}^{j=n} \langle A_j \left( \frac{\partial f}{\partial x_j} \right), f \rangle + \langle B(f), f \rangle$$

Since  $f_\nu(x, s)$  vanish when  $x \in \Omega \setminus K$  Stokes theorem gives

$$(ii) \quad 0 = \int \partial_{x_j} (\langle A_j(f), f \rangle) dx \quad : 1 \leq j \leq n$$

Rules for differentiation identifies for each  $j$  the integrand with

$$\left\langle \frac{\partial A_j}{\partial x_j}(f), f \right\rangle + \langle A_j(\partial_{x_j}(f)), f \rangle + \langle A_j(f), \partial_{x_j}(f) \rangle$$

Since  $A_j$  is hermitian we have

$$(iii) \quad \Re \langle A_j(\partial_{x_j}(f), f) = \Re \langle A_j(f), \partial_{x_j}(f) \rangle$$

Hence (ii) and (iii) give

$$\Re \int \langle A_j(\frac{\partial f}{\partial x_j}), f \rangle dx = -\frac{1}{2} \int \langle \frac{\partial A_j}{\partial x_j}(f), f \rangle dx$$

Introduce the matrix-valued function

$$\operatorname{div}(A) = \sum_{j=1}^{j=n} \frac{\partial A_j}{\partial x_j}$$

From the above we get the equation

$$(iv) \quad \frac{dJ}{ds} = \Re \int \left[ -\frac{1}{2} \langle \operatorname{div}(A)(f), f \rangle + \langle B(f), f \rangle \right] dx$$

By assumption the elements of the matrices  $\{A_j\}$  and of  $B$  as well as  $\{\frac{\partial A_j}{\partial x_j}\}$  are bounded  $C^\infty$ -functions. Hence (iv) and the Cauchy-Schwarz inequality gives a constant  $C$  such that the absolute value in the right hand side in (iv) is estimated above by

$$C \cdot \int |f(x, s)|^2 dx \quad : -s_* \leq s < s^*$$

It follows that

$$\left| \frac{dJ}{ds} \right| \leq C \cdot J(s) \quad : 0 \leq s \leq s^*$$

At the same time (i) means that  $J(0) = 0$ . Hence Picard's uniqueness theorem to be exposed in § XX implies that if  $J(s) = 0$  when  $-s_* < s < s_*$ , i.e.

$$(0.3.1) \quad J(x, s) = 0 \quad : (x, s) \in \{|x| \leq r\} \times [0, s^*]$$

So we have

$$\int_{\Omega} |f(x, s)|^2 dx = 0$$

and  $\{f_\nu\}$  are continuous functions they vanish identically in  $\Omega \times (-s_*, s_*)$  as requested.

**0.3 A semi-global uniqueness result.** Let  $\{A_j(x, s)\}$  be Hermitian matrices defined in the whole of  $\mathbf{R}^{n+1}$  whose elements are bounded  $C^\infty$ -functions. Similarly  $B(x, s)$  is defined in the whole of  $\mathbf{R}^{n+1}$ .

**0.3.1 Proposition.** *There exists a positive number  $\rho$  which only depends on the matrices above such that if  $R > 0$  and a vector valued  $C^\infty$ -function  $f(x, s)$  is defined in the ball  $\{|x|^2 + s^2 < R^2\}$  where it is a solution to (\*) and satisfies*

$$f(x, 0) = 0 \quad : |x| < R$$

*Then it follows that*

$$f(x, s) = 0 \quad : x^2 + s^2 < \rho \cdot R^2$$

*Proof.* Choose a test-function  $\psi(x)$  in  $\mathbf{R}^n$  such that  $\psi = 1$  when  $|x| \leq 1$  and vanishes when  $|x| > 3/2$  while the values stay in  $[0, 1]$ . Set

$$(i) \quad \mu = \max_{1 \leq j \leq m} \sup_{(x, s)} \|A_j(x, s)\|$$

where the supremum is taken over all  $(x, s)$  in  $\mathbf{R}^{n+1}$  and we have taken the Hilbert-Schmidt norms of the  $A$ -matrices. With  $\epsilon > 0$  we set  $\phi(x) = \psi(\epsilon \cdot x)$  and construct the following  $m \times m$ -matrices

$$H(x, s) = \sum_{j=1}^{j=n} \frac{\partial \phi}{\partial x_j}(x) \cdot A_j(x, s)$$

$$\widehat{A}_j(x, s) = \phi(x) \cdot A_j(x, \phi(x)s) \quad : \quad \widehat{B}(x, s) = \phi(x) \cdot B(x, \phi(x)s)$$

Put

$$F(x, s) = f(x, \phi(x)s)$$

Since  $0 \leq \phi(x) \leq 1$  hold for all  $x$  it follows that  $F$  is defined in  $\{|x|^2 + s^2 < R^2\}$  and the construction of  $\phi$  gives

$$(ii) \quad F(x, s) = f(x, s) \quad : \quad |x| < \epsilon^{-1}$$

Rules for differentiation show that  $F$  satisfies the system

$$(E - H(x, s))\partial_s(F) = \sum_{j=1}^{j=n} \widehat{A}_j(x, s)\partial_{x_j}(F) + \widehat{B}(x, s)F \quad : \quad x^2 + s^2 < R^2$$

where  $E$  is the identity operator.

**A choice of  $\epsilon$ .** The partial derivatives of the test-function  $\psi$  are bounded by some constant  $C$  and we set

$$\epsilon^* = \frac{1}{2n \cdot C\mu}$$

It follows that

$$\left| \frac{\partial \phi}{\partial x_j} \right| = \epsilon \cdot \left| \frac{\partial \psi}{\partial x_j} \right| \leq \frac{1}{2n\mu}$$

By (i) this entails that

$$(iii) \quad \sup_{(x,s)} \|H(x, s)\| \leq \frac{1}{2}$$

Next, the support of  $\phi$  is contained in the ball  $\{|x| \leq \frac{3}{2\epsilon^*}\}$  so the vanishing in (xx) entails that  $F(x, s) = 0$  when

$$\frac{3}{2\epsilon^*} \leq |x| < R$$

By the condition (xx)  $R \geq R_*$  entails that

$$\frac{3}{2\epsilon^*} = 3nC\mu = R_*/2 \leq R/2$$

So  $R \geq R_*$  implies that

$$(v) \quad F_\epsilon(x, s) = 0 \quad : \quad \frac{R}{2} \leq |x| < \sqrt{R^2 - s^2}$$

Now (iv) implies that the hermitian matrix  $E - H(x, s)$  is invertible and by (iii)  $F$  satisfies system as in (0.2). The vanishing in (v) therefore implies that  $F_\epsilon(x, s) = 0$  hold when  $x^2 + s^2 < R^2$ .

FINISH

### § 1: Symmetric hyperbolic systems.

The main result in this section appears in Theorem 1.xx. Before it can be announced we need several preliminaries. We will study periodic functions. Let  $n$  be a positive integer and consider the  $(n + 1)$ -dimensional torus  $T^{n+1}$  with variables  $(x, s) = (x_1, \dots, x_n, s)$ . Denote by  $C^\infty(T^{n+1})$  the space of complex-valued  $C^\infty$ -functions which are  $2\pi$ -periodic in all the variables. Passing to the  $L^2$ -norm the closure of these functions give the complex Hilbert space  $L^2(T^{n+1})$  whose vectors are complex-valued functions  $f(x, s)$  which are square integrable on the  $(n + 1)$ -dimensional  $2\pi$ -periodic torus. Next, to each multi-index  $\alpha = (\alpha_1, \dots, \alpha_{n+1})$  one associates the differential operator

$$\partial^\alpha = \partial_{x_1}^{\alpha_1} \dots \partial_{x_n}^{\alpha_n} \cdot \partial_s^{\alpha_{n+1}}$$

If  $k$  is a positive integer an inner product is defined on  $C^\infty(T^{n+1})$  by

$$(1.1) \quad \langle f, g \rangle_{(k)} = \sum_{|\alpha| \leq k} \int \partial^\alpha(f) \cdot \overline{\partial^\alpha(g)} dx ds$$

Passing to the closure we obtain a Hilbert space denoted by  $H^{(k)}$  whose elements are  $L^2$ -functions  $g(x, s)$  such that the distribution derivatives  $\partial^\alpha(g)$  are square integrable when  $|\alpha| \leq k$ . From § XX we recall:

**The Fourier-Sobolev Lemma.** *If  $k \geq 1$  every  $g \in H^{(k)}$  is a periodic function of class  $C^1$  at least on  $\mathbf{T}^{n+1}$ .*

More generally, if  $m \geq 2$  we consider vector-valued functions  $f = (f_1, \dots, f_m)$  and get the Hilbert space  $H^{(k)}[m]$  whose vectors are  $m$ -tuples of functions in  $H^{(k)}$ . With  $f = (f_1, \dots, f_m)$  and  $g = (g_1, \dots, g_m)$  the inner product is defined as in (1,1):

$$(1.2) \quad \langle f, g \rangle_{(k)} = \sum_{\nu=1}^m \sum_{|\alpha| \leq k} \int \partial^\alpha(f_\nu) \cdot \overline{\partial^\alpha(g_\nu)} dx ds$$

With  $m \geq 1$  we consider a matrix-valued functions  $\{A_j(x, s) \dots A_n(x, s)\}$  where each  $A_j(x, s)$  is an  $m \times m$ -matrix whose elements are periodic complex-valued  $C^\infty$ -functions on  $T^{n+1}$ . Let  $B(x, s)$  be another matrix-valued functions whose elements also are periodic and of class  $C^\infty$ . Set

$$(1.3) \quad P(x, s, \partial_x, \partial_s) = E_m \cdot \partial_s - \sum_{j=1}^{j=n} A_j(x, s) \cdot \partial_{x_j} + B(x, s)$$

This differential operator acts on vector-valued functions  $f$ . Identifying the space of vector-valued and periodic  $C^\infty$ -functions with a subspace of  $H^{(k)}[m]$  one has the linear map

$$P: f \rightarrow P(f)$$

from  $C^\infty[m]$  into  $H^{(k)}[m]$ . Keeping  $k$  and  $m$  fixed we denote this linear map by  $T_0$ . It means that  $T_0$  is densely defined linear operator on  $H^{(k)}[m]$  whose domain of definition  $\mathcal{D}(T_0) = C^\infty[m]$ . We have also the densely defined linear operator  $T_1$  where

$$\mathcal{D}(T_1) = \{f \in H^{(k)}[m] : P(f) \in H^{(k)}[m]\}$$

By the general result in § XX the graph of  $T_1$  taken in the product  $H^{(k)}[m] \times H^{(k)}[m]$  is closed. Next, since  $T_0$  is densely defined there exists the adjoint operator  $T_0^*$ . By definition  $\mathcal{D}(T_0^*)$  consists of vectors  $g \in H^{(k)}[m]$  for which there exists a constant  $C(g)$  such that

$$|\langle T_0(f), g \rangle| \leq C(g) \|f\|_k \quad : f \in \mathcal{D}(T_0)$$

and for such  $g$ -vectors we get a unique vector  $T_0^*(g)$  such that

$$\langle T_0(f), g \rangle = \langle f, T_0^*(g) \rangle$$



**1.4 The case when  $\{A_j\}$  are hermitian.** An  $m \times m$ -matrix  $A(x, s) = \{a_{\nu, \mu}(x, s)\}$  whose elements are periodic  $C^\infty$ -functions is hermitian if

$$a_{\mu, \nu}(x, s) = \overline{a_{\nu, \mu}(x, s)}$$

hold for all pairs  $1 \leq \nu, \mu \leq m$ .

**1.5 Proposition.** *If  $A_1, \dots, A_n$  are hermitian it follows that  $\mathcal{D}(T_0^*) = \mathcal{D}(T_1)$  and there exists a bounded and self-adjoint linear operator  $\mathcal{B}$  on  $H^{(k)}[m]$  such that*

$$(*) \quad T_0^* + T_1 = \mathcal{B}$$

Before we enter the proof we need some constructions. Repeated use of Stokes Theorem gives the equality below for every pair of functions  $f, g$  in  $C^\infty(T^{n+1})$  and every multi-index  $\alpha$ :

$$(-1)^{|\alpha|} \cdot \int \partial^{2\alpha}(f) \cdot \bar{g} \, dx ds = \int \partial^\alpha(f) \cdot \overline{\partial^\alpha(g)} \, dx ds$$

More generally, let  $Q = Q(x, s, \partial_x, \partial_s)$  be a differential operator. With  $Q$  given as

$$Q = \sum q_\alpha(x, s) \cdot \partial^\alpha$$

where  $q_\alpha \in C^\infty(T^{n+1})$  one gets the differential operator

$$Q^* = \sum (-1)^\alpha \cdot \partial^\alpha \circ q_\alpha(x, s)$$

where  $\partial^\alpha \circ q_\alpha(x, s)$  is the product taken in the ring of differential operators with  $\{q_\alpha(x, s)\}$  regarded as zero-order differential operators. Stokes theorem gives

$$(1.6) \quad \int Q(f) \cdot \bar{g} \, dx ds = \int f \cdot Q^*(\bar{g}) \, dx ds$$

Let us write out

$$Q^* = \sum r_\alpha(x, s) \cdot \partial^\alpha$$

We take the complex conjugates of the  $r$ -functions and put

$$\overline{Q^*} = \sum \bar{r}_\alpha(x, s) \cdot \partial^\alpha$$

Using the hermitian inner product on  $L^2(T^{n+1})$  we can express (1.6) by the equation

$$(1.7) \quad \langle Q(f), g \rangle = \langle f, \overline{Q^*}(g) \rangle$$

**1.8 The  $\Gamma$ -operator.** Let us introduce the differential operator

$$\Gamma = \sum_{|\alpha| \leq k} (-1)^{|\alpha|} \cdot \partial^{2\alpha}$$

If  $m \geq 2$  we denote by  $\Gamma_m$  the operator given by the diagonal  $m \times m$ -matrix where whose diagonal elements are  $\Gamma$ . Stokes theorem entails that if  $f$  and  $g$  is a pair of vector-valued functions in  $C^\infty(T^{n+1})$  then

$$(1.9) \quad \langle f, g \rangle_{(k)} = \int \Gamma_m(f) \cdot \bar{g} \, dx ds$$

**1.10 Exercise.** Both sides in (1.9) are defined under the relaxed condition that  $g \in H^{(k)}[m]$  while  $f \in C^\infty[m]$ . Show by continuity that (1.9) remains valid for such pairs.

Next, in the algebra of  $m \times m$ -matrices whose elements are differential operators we consider the product  $\Gamma_m \cdot P$ .

**Exercise.** Show that when  $P$  is as in (1.3) where  $\{A_j\}$  are hermitian then there exists an  $m \times m$ -matrix  $Q$  whose elements are differential operators of order  $\leq 2k$  such that the following hold in the algebra above:

$$(1.11) \quad \Gamma_m \circ P + \overline{P^*} \circ \Gamma_m = Q$$

Now (xx) and (xx) give the equation

$$\langle T_0(f), g \rangle_{(k)} = \int Q(f) \cdot \bar{g} \, dx ds - \int \overline{P^*} \circ \Gamma_m(f) \cdot \bar{g} \, dx ds$$

Apply (xx) to the pair of vector-valued functions  $\Gamma_m(f)$  and  $g$  and the differential operator  $\overline{P^*}$ . Notice that the complex conjugate of the adjoint  $(\overline{P^*})^*$  is equal to  $P$  and from this the reader can check from the above that the last term in (xx) is equal to

$$- \int \Gamma_m(f) \cdot \overline{P(g)} \, dx ds$$

Applying (xx) this entails that the following hold for each pair  $f, g$  in  $C^\infty(T^{n+1})$ .

$$\langle T_0(f), g \rangle_{(k)} = -\langle f, T_0(g) \rangle_{(k)} + \int Q(f) \cdot \bar{g} \, dx ds$$

**Exercise.** Since the differential operator  $Q$  has degree  $\leq 2k$  the reader should verify the existence a bounded linear operator  $\mathcal{B}_k$  on the Hilbert space  $H^{(k)}[m]$  such that

$$\int Q(f) \cdot \bar{g} \, dx ds = \langle \mathcal{B}_k(f), g \rangle_{(k)}$$

hold when  $f$  is a vector-valued  $C^\infty$ -function and  $g \in H^{(k)}[m]$ . In particular we can take a pair  $f, g$  in  $C^\infty$  and notice that

$$(f, g) \mapsto \langle T_0(f), g \rangle_{(k)} + \langle T_0(g), f \rangle_{(k)}$$

is symmetric in  $f$  and  $g$ . From this the reader can conclude that the bounded operator  $\mathcal{B}_k$  is self-adjoint.

*Proof of Proposition 1.5*

Assume first that  $g \in \mathcal{D}(T_0)$  which gives

$$\langle T_0(f), g \rangle_{(k)} = \langle f, T_0^*(g) \rangle_{(k)}$$

Here  $g \in H^{(k)}$  and regarding  $g$  as a distribution we get the vector-valued distribution  $P(g)$ . Now xx hold for all vector-valued periodic  $C^\infty$ -functions  $f$ . From the above the trinsgle inequality and Cauchy-Schwarz gives

$$|\langle f, P(f)g \rangle_{(k)}| \leq (\|g\|_k + \|\mathcal{B}_k(g)\|_k) \cdot \|f\|_k$$

Since this inequqlity hold for all  $f$  in the dense subspace  $C^\infty[m]$  it follows that the distribution  $P(f)g$  belongs to  $H^{(k)}[m]$  so by the construction of  $T_1$  one has  $g \in \mathcal{D}(T_1)$ . Hence one has the inclusion

$$(i) \quad \mathcal{D}(T_1) \subset \mathcal{D}(T_0^*)$$

Conveersley, if  $g \in \mathcal{D}(T_1)$  the absolute value in the right hand side of (xx) is majorized by

$$\|T_1(g)\|_k + \|\mathcal{B}_k(g)\|_k \cdot \|f\|_k$$

The construction of  $T_0^*$  entails that  $g \in \mathcal{D}(T_0^*)$  and hence equality holds in (i) Finally it is clear that this equality and (xxx) gives the operator equation

$$T_0^* = -T_1 + \mathcal{B}_k^*$$

Since we already proved that  $\mathcal{B}_k$  is self-adjoint the proof of Proposition 1.5 is finished.

## § 2. A study of $T_1$

In § 1 we constructed the densely defined and closed operator  $T_1$  on  $H^{(k)}[m]$ . Consider some  $f \in C^\infty[m]$  and a real number  $\lambda$ . Now

$$\|T_1(f) + \lambda \cdot f - \frac{1}{2}\mathcal{B}_k^*(f)\|_{(k)}^2 =$$

$$\|T_1(f) - \frac{1}{2}\mathcal{B}_k^*(f)\|_{(k)}^2 + \lambda^2 \|f\|_{(k)}^2 + \lambda \cdot \langle T_1(f) - \frac{1}{2}\mathcal{B}_k^*(f), f \rangle_{(k)} + \lambda \cdot \langle f, T_1(f) - \frac{1}{2}\mathcal{B}_k^*(f) \rangle_{(k)}$$

Since  $f$  is  $C^\infty$  we have  $T_1(f) = T_0(f)$  and since  $\mathcal{B}_k^*$  is self-adjoint it follows that the sum of the last two terms above becomes

$$(i) \quad \lambda \cdot (\langle f, T_0^*(f) - \frac{1}{2}\mathcal{B}_k^*(f) \rangle_{(k)} + \langle f, T_0(f) - \frac{1}{2}\mathcal{B}_k^*(f) \rangle_{(k)})$$

The operator equation in Proposition 1.5 gives

$$T_0(f) + T_0^*(f) = \mathcal{B}_k^*$$

which proves that (i) is zero. Hence we have proved the equality

$$(2.1) \quad \|T_1(f) + \lambda \cdot f - \frac{1}{2}\mathcal{B}_k^*(f)\|_{(k)}^2 = \|T_1(f) - \frac{1}{2}\mathcal{B}_k^*(f)\|_{(k)}^2 + \lambda^2 \|f\|_{(k)}^2$$

Since  $\|T_1(f) - \frac{1}{2}\mathcal{B}_k^*(f)\|_{(k)} \geq 0$  the triangle inequality gives the inequality below for every real number  $\lambda$

$$(2.2) \quad \|T_1(f) + \lambda \cdot f\|_{(k)} \geq |\lambda| \cdot \|f\|_{(k)} - \frac{1}{2} \cdot \|\mathcal{B}_k^*(f)\|_{(k)} \quad : f \in C_{\text{per}}^\infty[m]$$

Above the real number  $\lambda$  can be both positive or negative and the inequality is of interest when  $|\lambda|$  exceeds the operator norm of  $\frac{\mathcal{B}_k^*}{2}$ . We have for example

$$(2.3) \quad \|T_1(f) + \lambda \cdot f\|_{(k)} \geq \frac{|\lambda|}{2} \cdot \|f\|_{(k)} \quad : |\lambda| \geq \|\mathcal{B}_k^*\|$$

**2.4 The operator  $\widehat{T}_0$ .** Recall that  $T_1$  is closed and extends  $T_0$  in the sense that its graph contains that of  $T_0$ . Taking the closure of  $\Gamma(T_0)$  we get the densely defined and closed operator  $\widehat{T}_0$  whose graph is contained in  $\Gamma(T_1)$ . When  $f$  are  $C^\infty$ -functions we have  $T_0(f) = \widehat{T}_0(f) = T_1(f)$  so (2.3) holds with  $T_1$  replaced by  $\widehat{T}_0$ . Since  $C^\infty[m]$  is dense in  $H^{(k)}[m]$  the inequality below holds for every  $g \in \mathcal{D}(\widehat{T}_0)$ :

$$(2.4) \quad \|\widehat{T}_0(g) + \lambda \cdot g\|_{(k)} \geq \frac{|\lambda|}{2} \cdot \|g\|_{(k)} \quad : |\lambda| \geq \|\mathcal{B}_k^*\|$$

Since  $\widehat{T}_0$  is closed it follows that the range of  $\widehat{T}_0(g) + \lambda \cdot E$  is closed when  $|\lambda| \geq \|\mathcal{B}_k^*\|$ .

**2.5 Density of the range.** From now on  $|\lambda| \geq \|\mathcal{B}_k^*\|$ . Recall from the general material in § XX that the adjoint of  $\widehat{T}_0$  is equal to  $T_0^*$ . If the range of  $\widehat{T}_0(g) + \lambda \cdot E$  is not dense there exists  $0 \neq g \in H^{(k)}[m]$  such that

$$\langle \widehat{T}_0(f) + \lambda \cdot f, g \rangle_{(k)} = 0 \quad : f \in C^\infty[m]$$

Here  $\widehat{T}_0(f) = P(f)$  for  $C^\infty$ -functions and (x) gives

$$|\langle P(f), g \rangle_{(k)}| \leq |\lambda| \cdot |\langle f, g \rangle_{(k)}|$$

It follows that  $g \in \mathcal{D}(T_0^*)$  so (i) in gives

$$\langle f, T_0^*(g) \rangle_{(k)} + \lambda \cdot \langle f, g \rangle_{(k)} = 0$$

This hold for all  $f \in C^\infty[m]$  and since  $\lambda$  is real we get hence

$$T_0^*(g) + \lambda \cdot g = 0$$

Now the operator equation in Proposition XX gives

$$T_1(g) = \lambda \cdot g + \mathcal{B}_k^*(g)$$

At this stage we assume that  $k$  is so large that the Sobolev inequality entails that  $H^{(k)}[m]$  consists of vector-valued functions of class  $C^1$  at least which in addition are peridic on the whole torus  $T^{n+1}$ . Now (xx) means that

$$P(g) = \lambda \cdot g + \mathcal{B}_k^*(g)$$

From this we shall prove that  $g = 0$  if the real number  $\lambda$  is sufficiently large. To attain this we consider the function

$$G(s) = \int_{T^n} |g(x, s)|^2 dx$$

It follows that

$$\frac{dG}{ds} = 2 \cdot \Re \int_{T^n} \partial_s(g)(x, s) \cdot \overline{g(x, s)} dx$$

Then conclude...

**2.6 Conclusive results.** If  $k \geq xxx$  we have proved that there exists a positive constant  $\mu_k$  such that if  $|\lambda| \geq \mu_k$  then the densely defined operator  $\lambda \cdot E - \widehat{T}_0$  is surjective and at the same time one has the inequality XX. This means that there exists the resolvent operator  $R(\lambda; \widehat{T}_0)$  for such real  $\lambda$ . Keeping  $k$  fixed we get the closed spectrum of  $\widehat{T}_0$  which by the general result in § xx is a closed subset of  $\mathbf{C}$ . Since  $\widehat{T}_0$  is an unbounded operator one cannot expect that the spectrum is compact. Moreover in contrast to the more favourable cases for elliptic equations the resolvent operators are in general not compact.

### § 1. Differential inequalities.

Let  $M(s)$  be a non-negative real-valued continuous function on a closed interval  $[0, s^*]$ . To each  $0 \leq s < s^*$  we set

$$d_M^+(s) = \limsup_{\Delta s \rightarrow 0} \frac{M(s + \Delta s) - M(s)}{\Delta s}$$

where  $\Delta s$  are positive during the limit.

**1.1 Proposition.** *If there exists a real number  $B$  such that  $d_M^+(s) \leq B \cdot M(s)$  holds in  $[0, s^*)$  then*

$$M(s) \leq M(0) \cdot e^{Bs} \quad : 0 < s \leq s^*$$

The proof of this result is left as an exercise to the reader. The hint is to consider the function  $N(s) = M(s)e^{-Bs}$  and show that  $d_N^+(s) \leq 0$  for all  $s$ . Notice that  $B$  is an arbitrary real number, i.e. it may also be  $< 0$ .

More generally, let  $k(s)$  be some non-decreasing continuous function with  $k(0) = 0$ . suppose that

$$d_M^+(s) \leq B \cdot M(s) + k(s) \quad : 0 \leq s < s^*$$

Now the reader may verify that

$$(1.1.1) \quad M(s) \leq M(0) \cdot e^{Bs} + \int_0^s k(t) dt$$

Next, consider a product set  $\square = [0, \pi] \times [0, s^*]$  where  $0 \leq x \leq \pi$  and consider functions  $g(x, s)$  which are periodic in  $x$ , i.e.

$$g(0, s) = g(\pi, s) \quad : 0 \leq s \leq s^*$$

A  $C^1$ -function  $g$  is periodic  $C^1$ -function when  $g$  and the partial derivatives  $\partial_s(g)$  and  $\partial_x(g)$  are periodic in  $x$ .

**1.2 Theorem.** *Let  $g$  be a  $C^1$ -function which satisfies the PDE-equation*

$$(*) \quad \partial_s(g)(x, s) = a \cdot \partial_x(g) + b \cdot g$$

in  $\square$  where  $a$  and  $b$  are  $x$ -periodic real-valued continuous functions. Set

$$M_g(s) = \max_x |g(x, s)| \quad : B = \max_{x,s} |b(x, s)|$$

Then one has the inequality

$$M_g(s) \leq M_g(0) \cdot e^{Bs}$$

*Proof.* Consider some  $0 < s < s^*$  and let  $\epsilon > 0$ . Put

$$m^*(s) = \{x : g(x, s) = M_g(s)\}$$

The continuity of  $g$  entails that the function  $M(s)$  is continuous and the sets  $m^*(s)$  are compact. If  $x^* \in m^*(s)$  the periodicity of the  $C^1$ -function  $x \mapsto g(x, s)$  entails that  $\partial_x(x^*, s) = 0$  and (\*) gives

$$\partial_s(g)(x, s) = b(x, s)g(x, s) \quad : x \in m^*(s)$$

Next, let  $\epsilon > 0$ . We find an open neighborhood  $U$  of  $m^*(s)$  such that

$$|\partial_x(g)(x, s)| \leq \epsilon \quad : x \in U$$

Now there exists  $\delta > 0$  such that

$$|g(x, s)| \leq M(s) - 2\delta \quad : x \in [0, \pi] \setminus U$$

Continuity gives some  $\rho > 0$  such that if  $0 < \Delta s < \rho$  then the inequalities below hold:

$$(i) \quad |g(x, s + \Delta s)| \leq M(s) - \delta \quad : x \in [0, \pi] \setminus U \quad : M(s + \Delta s) > M(s) - \delta$$

$$(ii) \quad M(s + \Delta s) \leq M(s) + \epsilon \quad : \quad |\partial_x(g)(x, s + \Delta s)| \leq 2\epsilon \quad : \quad x \in m^*(s)$$

If  $0 < \Delta s < \rho$  we see that (i) gives  $x \in m^*(s + \Delta s) \subset U$  and for such  $x$ -values Rolle's mean-value theorem and the PDE-equation give

$$M_g(x, s + \Delta s) - g(x, s) = \Delta s \cdot \partial_s(g(x, s + \theta \cdot \Delta s)) =$$

$$(iii) \quad \Delta s \cdot [a(x, s + \Delta s) \cdot \partial_x(g)(x + \theta \cdot \Delta s) + b(x, s + \Delta s) \cdot g(x, s + \theta \cdot \Delta s)]$$

Let  $A$  be the maximum norm of  $|a(x, s)|$  taken over  $\square$ . Since  $|g(x, s)| \leq M(s)$  the triangle inequality and (iii) give

$$M(s + \Delta s) \leq M(s) + \Delta s[A \cdot 2\epsilon + B \cdot M(s + \theta \cdot \Delta s)]$$

Since the function  $s \mapsto M(s)$  is continuous it follows that

$$\limsup_{\Delta s \rightarrow 0} \frac{M(s + \Delta s) - M(s)}{\Delta s} \leq A \cdot 2\epsilon + BM(s)$$

Above  $\epsilon$  can be arbitrary small and hence

$$d^+(s) \leq B \cdot M(s)$$

Then Proposition 1.1 gives (\*) in the theorem.

**1.3 Higher order derivatives.** Suppose that  $g$  is a  $C^2$ -function satisfying the PDE-equation (\*) where  $a$  and  $b$  have a continuous partial  $x$ -derivative. Set  $h = \partial_x(g)$ . Since the differential operators  $\partial_x$  and  $\partial_s$  commute we obtain

$$(1.3.1) \quad \partial_s(h) = \partial_x(a \cdot h) + \partial_x(b \cdot g) = a \cdot \partial_x h + (\partial_x(a) + b)h + \partial_x(b)g$$

**$L^2$ -inequalities.** Let  $g(x, s)$  be a  $C^1$ -function satisfying (\*) in Theorem 1.2. Set

$$J_g(s) = \int_0^\pi g^2(x, s) dx$$

Differentiation with respect to  $s$  and (\*) give

$$\frac{dJ_g}{ds} = 2 \cdot \int_0^\pi (a\partial_s(g) \cdot \partial g + b \cdot g) dx$$

By periodicity  $\int_0^\pi \partial_x(ag^2) dx = 0$  which entails that the right hand side becomes

$$\int_0^\pi (-\partial_x(a) + b) \cdot g^2 dx$$

So if  $K$  is the maximum norm of  $-\partial_x(a) + b$  over  $\square$  it follows that

$$\frac{dJ_g}{ds}(s) \leq K \cdot J(s)$$

Hence Theorem xx gives

$$\int_0^\pi g^2(x, s) dx \leq e^{Ks} \cdot \int_0^\pi g^2(x, 0) dx \quad : \quad 0 < s \leq s^*$$

Integration with respect to  $s$  entails that

$$\iint_{\square} g^2(x, s) dx ds \leq \int_0^{s^*} e^{Ks} ds \cdot \int_0^\pi g^2(x, 0) dx$$

This, the  $L^2$ -integral of  $x \rightarrow g(x, 0)$  majorizes both the area integral and each slice integral when  $0 < s \leq s^*$ .

**Higher order derivatives.** Same procedure gives majorisations of these integrals when higher order  $x$ -derivatives are inserted. Simialrly we regard PDE-equations when  $g$  is replaced by  $h = \partial_s(g)$  and so on.

## § 2. A boundary value equation

Let  $a(x, s)$  and  $b(x, s)$  be real-valued  $C^\infty$ -functions on  $\square$  which are periodic in  $x$ . Consider the PDE-operator

$$P = \partial_s - a \cdot \partial_x - b$$

Given a periodic  $C^1$ -function  $f(x)$  on  $[0, \pi]$  we seek a  $C^1$ -function  $g(x, s)$  in  $\square$  which satisfies  $P(g) = 0$  and the initial condition

$$g(x, 0) = f(x)$$

From § xx we see that  $g$  is unique if it exists. There remains to prove existence.

**2.1 Theorem.** *For every positive integer  $p$  and each periodic  $f \in C^p[0, \pi]$  there exists a unique periodic  $g \in C^p(\square)$  where  $P(g) = 0$  and  $g(x, 0) = f(x)$ .*

The proof requires several steps and employs Hilbert space methods. To each non-negative integer  $k$  we get the complex Hilbert space  $\mathcal{H}^{(k)}$  from § xx, i.e. we complete the space of complex-valued  $C^k$ -functions on  $\square$  which are periodic with respect to  $x$ . Sobolev's inequality shows that if  $k \geq 2$  then every function in  $\mathcal{H}^{(k)}$  is continuous and more generally one has the inclusion

$$\mathcal{H}^{(k)} \subset C^{k-2}(\square) \quad : k \geq 3$$

Moreover, the first order PDE-operator  $P$  maps  $\mathcal{H}^{(k+1)}$  into  $\mathcal{H}^{(k)}$ . From now on we only consider  $k$ -integers which are  $\geq 2$ . On the periodic  $x$ -interval  $[0, \pi]$  we get the Hilbert spaces  $\mathcal{H}^k[0, \pi]$ . The result in § xx shows that if  $f(x) \in \mathcal{H}^k[0, \pi]$  is such that there exists some  $F(x, s) \in \mathcal{H}^{(k)}$  such that  $P(F) = 0$  and  $F(x, 0) = f(x)$  then  $F$  is unique and there exists a constant  $C$  which only depends upon the  $C^\infty$ -functions  $a$  and  $b$  such that

$$(i) \quad \|F\|_k \leq C \cdot \|f\|_k$$

where we have taken norms in  $\mathcal{H}^{(k)}$   $\mathcal{H}^k[0, \pi]$ . Moreover  $C$  can be chosen such that the function  $f^*(x) = F(x, s^*)$  satisfies

$$\|f^*\|_k \leq C \cdot \|f\|_k$$

Let  $\mathcal{D}_k(P)$  be the set of all  $f(x) \in \mathcal{H}^k[0, \pi]$  for which  $F(x, s)$  above exists.

**2.2 Density Lemma.** *If  $\mathcal{D}_k(P)$  is dense in  $\mathcal{H}^k[0, \pi]$  then it is equal to  $\mathcal{H}^k[0, \pi]$ .*

*Proof.* Let  $f$  be in  $\mathcal{H}^k[0, \pi]$  and by density we find a sequence  $\{f_n\}$  in  $\mathcal{D}_k(P)$  where  $\|f_n - f\|_k \rightarrow 0$ . By (i) we have

$$\|F_n - F_m\|_k \leq C \|f_n - f_m\|_k$$

Hence  $\{F_n\}$  is a Cauchy sequence in the Hilbert space  $\mathcal{H}^{(k)}$  and converges to a limit  $F$ . Since each  $P(F_n) = 0$  it follows that  $P(F) = 0$  and it is clear that the continuous boundary value function  $F(x, 0) = f(x)$  which gives  $f \in \mathcal{D}_k(P)$ .

**2.3 The operators  $S_k$ .** To each  $f \in \mathcal{D}_k(P)$  we get  $F(x, \pi)$  in  $\mathcal{H}^k[0, \pi]$  and set

$$S_k(f) = F(x, \pi)$$

So the domain of definition of  $S_k$  is  $\mathcal{D}_k(P)$  and (ii) gives a constant  $M_k$  such that

$$\|S_k(f)\| \leq M_k \cdot \|f\|_k \quad : f \in \mathcal{D}_k(P)$$

**2.4 Proposition.** *For each  $k$  there exists some  $\alpha(k) > 0$  such that the range of the operator  $E - \alpha \cdot S_k$  contains all periodic  $C^\infty$ -functions on  $[0, \pi]$ .*

We prove Proposition 2.4 in § XXX. Let us now show that it gives the density of  $\mathcal{D}_k(P)$ . Namely, if  $\mathcal{D}_k(P)$  fails to be dense there exists a non-zero  $f_0 \in \mathcal{D}_k(P)$  which

is  $\perp$  to  $\mathcal{D}_k(P)$  and normalised so that  $\|f_0\|_k = 1$ . In Proposition 2.4 we choose  $0 < \alpha \leq \alpha(k)$  so small that

$$(i) \quad \alpha < \frac{1}{2M_k}$$

Since periodic  $C^\infty$ -functions are dense in  $\mathcal{H}^k[0, \pi]$  it follows from Proposition 2.4 that there exists a sequence  $\{h_n\}$  in  $\mathcal{D}_k(P)$  such that

$$(ii) \quad \lim_{n \rightarrow \infty} \|h_n - \alpha \cdot S_k(h_n) - f_0\|_k \rightarrow 0 \implies$$

$$(iii) \quad \langle f_0, f_0 \rangle = 1 = \lim \langle f_0, h_n - \alpha \cdot S_k(h_n) \rangle = -\alpha \cdot \lim \langle f_0, S_k(h_n) \rangle$$

Next, the triangle inequality and (ii) give

$$(iv) \quad \|h_n\|_k \leq 1 + \alpha \cdot \|S_k(h_n)\| \leq 1 + 1/2 \cdot \|h_n\| \implies \|h_n\|_k \leq 2$$

Finally, by the Cauchy-Schwarz inequality the absolute value in the right hand side of (iii) is majorized by

$$\alpha \cdot M_k \cdot 2 < 1$$

which contradicts (iii). Hence the orthogonal complement of  $\mathcal{D}_k(P)$  is zero which proves the requested density and by the above we get the following conclusive result:

**2.5 Theorem.** *If  $k \geq 2$  and  $f(x) \in \mathcal{H}^k[0, \pi]$  there exists a unique  $F(x, s) \in \mathcal{H}^{(k)}$  such that  $F(x, 0) = f(x)$ .*



### § 3. A class of inhomogeneous PDE-equations.

Before Theorem 3.1 is announced we introduce some notations. Put

$$\square = \{0 \leq x \leq \pi\} \times \{0 \leq s \leq 2\pi\}$$

We shall consider doubly periodic functions  $g(x, s)$  on  $\square$ , i.e.

$$g(\pi, s) = g(0, s) \quad : \quad g(x, 0) = g(x, 2\pi)$$

If  $k \geq 0$  we denote by  $C^k(\square)$  the space of  $k$ -times doubly-periodic continuously differentiable functions. If  $g \in C^k(\square)$  we set

$$\|g\|_{(k)}^2 = \sum_{j,\nu} \int_{\square} \left| \frac{\partial^{j+\nu} g}{\partial x^j \partial s^\nu}(x, s) \right|^2 dx ds$$

with the double sum extended pairs  $j + \nu \leq k$ . This gives the complex Hilbert space  $\mathcal{H}^{(k)}$  after a completion of  $C^k(\square)$  with respect to the norm above. Recall from § xx that every function  $g \in \mathcal{H}^{(2)}$  is automatically continuous and doubly periodic on the closed square. More generally, if  $k \geq 3$  each  $g \in \mathcal{H}^{(k)}$  has continuous and doubly periodic derivatives up to order  $k - 2$ . Next, consider a first order PDE-operator

$$P = \partial_s - a(x, s)\partial_x - b(x, s)$$

where  $a$  and  $b$  are real-valued doubly periodic  $C^\infty$ -functions. It is clear that  $P$  maps  $\mathcal{H}^{(k)}$  into  $\mathcal{H}^{(k+1)}$  for every  $k \geq 2$ . Keeping  $k \geq 2$  fixed we set

$$\mathcal{D}_k(P) = \{g \in \mathcal{H}^{(k)} : P(g) \in \mathcal{H}^{(k)}\}$$

Since  $C^\infty(\square)$  is dense in  $\mathcal{H}^{(k)}$  this yields for each  $k \geq 2$  a densely defined operator

$$(i) \quad P: \mathcal{D}_k(P) \rightarrow \mathcal{H}^{(k)}$$

In  $\mathcal{H}^{(k)} \times \mathcal{H}^{(k)}$  we get the graph

$$\Gamma_k = \{(g, P(g)) : g \in \mathcal{D}_k(P)\}$$

Since  $P$  is a differential operator the general result in § xx entails that  $\Gamma_k$  is a closed subspace so the densely defined operator in (i) has a closed graph. Thus, for each  $k \geq 2$  we have a densely defined linear operator and closed operator on  $\mathcal{H}^{(k)}$  denoted by  $\mathcal{T}_k$ . So its domain of definition  $\mathcal{D}(\mathcal{T}_k) = \mathcal{D}_k$ . Next, we consider the graph

$$\gamma_k = \{(g, P(g)) : g \in C^\infty(\square)\}$$

This is a subspace of  $\Gamma_k$  and the closure  $\bar{\gamma}_k$  yields the graph of another densely defined linear operator denoted by  $T_k$ . We remark that the inclusion

$$\mathcal{D}(T_k) \subset \mathcal{D}(\mathcal{T}_k)$$

in general is strict. Let  $E$  be the identity operator. With these notations one has

**3.1 Theorem.** *For each integer  $k \geq 2$  there exists a positive real number  $\rho(k)$  such that  $T_k - \lambda \cdot E$  is surjective on  $\mathcal{H}^{(k)}$  for every  $\lambda > \rho(k)$  and its kernel is zero.*

**3.2 Remark.** The result is remarkable since  $T_k$  is only densely defined while Theorem 3.1 asserts that

$$T_k - \lambda \cdot E: \mathcal{D}_k \rightarrow \mathcal{H}^{(k)}$$

is bijective. Hence the closed and densely defined operator  $T_k$  has a non-empty resolvent set so by the general results in §§ x there exists resolvent operators  $R_k(\lambda)$  defined outside the closed spectrum  $\sigma(T_k)$  where the composed operators

$$(\lambda \cdot E - T_k) \circ R_k(\lambda) = E$$

for all  $\lambda$  outside  $\sigma(T_k)$ . The determination of these spectra is unclear and most likely one needs extensive numerical studies to grasp these closed sets which in addition depend upon  $k$ .

Next, recall from § xx that the closed and densely defined operator  $T_k$  has an adjoint  $T_k^*$ . A crucial step in the proof of Theorem 3.1 is the following:

**3.3 Proposition.** *One has the equality  $\mathcal{D}(T_k^*) = \mathcal{D}_k$  and there exists a bounded self-adjoint operator  $B_k$  on  $\mathcal{H}^{(k)}$  such that*

$$T_k^* = -\mathcal{T}_k + B_k$$

*Proof of Proposition 3.3*

Keeping  $k \geq 2$  fixed we set  $\mathcal{H} = \mathcal{H}^{(k)}$ . For each pair  $g, f$  in  $\mathcal{H}$  their inner product is defined by

$$\langle f, g \rangle = \sum \int_{\square} \frac{\partial^{j+\nu} f}{\partial x^j \partial s^\nu}(x, s) \cdot \overline{\frac{\partial^{j+\nu} g}{\partial x^j \partial s^\nu}(x, s)} dx ds$$

where the sum is taken when  $j + \nu \leq k$ . Introduce the differential operator

$$\Gamma = \sum_{j+\nu \leq k} (-1)^{j+\nu} \cdot \partial_x^{2j} \cdot \partial_s^{2\nu}$$

Partial integration gives

$$(i) \quad \langle f, g \rangle = \int_{\square} f \cdot \Gamma(\bar{g}) dx ds = \int_{\square} \Gamma(f) \cdot \bar{g} dx ds \quad : f, g \in C^\infty$$

Now we consider the operator  $P = \partial_s - a \cdot \partial_x - b$  and get

$$(ii) \quad \langle P(f), g \rangle = \int_{\square} P(f) \cdot \Gamma(\bar{g}) dx ds$$

Partial integration identifies (ii) with

$$(iii) \quad - \int_{\square} f \cdot (\partial_s - \partial_x(a) - a \cdot \partial_x - b) \circ \Gamma(\bar{g}) dx ds$$

**1.1 Exercise.** In (iii) appears the composed differential operator

$$\partial_s - \partial_x(a) - a \cdot \partial_x - b) \circ \Gamma$$

Show that in the ring of differential operators with  $C^\infty$ -coefficients this differential operator can be written in the form

$$\Gamma \circ (\partial_s - a \cdot \partial_x - b) + Q(x, s, \partial_x, \partial_s)$$

where  $Q$  is a differential of order  $\leq 2k$  with coefficients in  $C^\infty(\square)$ . Conclude from the above that

$$(1.1.1) \quad \langle Pf, g \rangle = -\langle f, Pg \rangle + \int_{\square} f \cdot Q(\bar{g}) dx ds$$

**1.2 Exercise.** With  $Q$  as above we have a bilinear form which sends a pair  $f, g$  in  $C^\infty(\square)$  to

$$(1.2.1) \quad \int_{\square} f \cdot Q(\bar{g}) dx ds$$

Use partial integration and the Cauchy-Schwarz inequality to show that there exists a constant  $C$  which depends on  $Q$  only such that the absolute value of (1.2.1) is majorized by  $C_Q \cdot \|f\|_k \cdot \|g\|_k$ . Conclude that there exists a bounded linear operator  $B_k$  on  $\mathcal{H}$  such that

$$(1.2.2) \quad \langle f, B_k(g) \rangle = \int_{\square} f \cdot Q(\bar{g}) dx ds$$

**1.3 Proof that  $B_k$  is self-adjoint** From the above we have

$$(1.3.1) \quad \langle Pf, g \rangle = -\langle f, Pg \rangle + \langle f, B_k(g) \rangle$$

Keeping  $f$  in  $C^\infty(\square)$  we notice that  $\langle f, B_k(g) \rangle$  is defined for every  $g \in \mathcal{H}$ . From this the reader can check that (1.3.1) remains valid when  $g$  belongs to  $\mathcal{D}(\mathcal{T}_k)$  which means that

$$(1.3.2) \quad \langle Pf, g \rangle = -\langle f, \mathcal{T}_k g \rangle + \langle f, B_k(g) \rangle \quad : f \in C^\infty(\square)$$

Moreover, when both  $f$  and  $g$  belong to  $C^\infty(\square)$  we can reverse their positions in (\*) which gives

$$(1.3.3) \quad \langle Pg, f \rangle = -\langle g, Pf \rangle + \langle g, B_k(f) \rangle$$

Since  $a$  and  $b$  are real-valued it is clear that

$$(1.3.4) \quad \langle Pg, f \rangle = -\langle f, Pg \rangle$$

It follows that

$$(1.3.5) \quad \langle f, B_k(g) \rangle = \langle g, B_k(f) \rangle \quad : f, g \in C^\infty(\square)$$

Since this hold for all pairs of  $C^\infty$ -functions and  $B_k$  is a bounded linear operator on  $\mathcal{H}$  the density of  $C^\infty(\square)$  entails that  $B_k$  is a bounded self-adjoint operator on  $\mathcal{H}$ .

**1.4 The equality  $\mathcal{D}(T_k^*) = \mathcal{D}_k$ .** The density of  $C^\infty(\square)$  in  $\mathcal{H}$  entails that a function  $g \in \mathcal{H}$  belongs to  $\mathcal{D}(T_k^*)$  if and only if there exists a constant  $C$  such that

$$(1.4.1) \quad |\langle Pf, g \rangle| \leq C \cdot \|f\| \quad : f \in C^\infty(\square)$$

Since  $B_k$  is a bounded operator, (1.3.2) gives the inclusion

$$(1.3.3) \quad \mathcal{D}_k \subset \mathcal{D}(T_k^*)$$

To prove the opposite inclusion we use that the  $\Gamma$ -operator is elliptic. If  $g \in \mathcal{D}(T_k^*)$  we have from (i) in § 1.1:

$$\langle Pf, g \rangle = \langle f, T_k^* g \rangle = \int \Gamma(f) \cdot \overline{T_k^*(g)} dx ds \quad : f \in C^\infty(\square)$$

Similarly

$$\langle f, B_k(g) \rangle = \int \Gamma(f) \cdot \overline{B_k(g)} dx ds$$

Treating  $\mathcal{T}_k(g)$  as a distribution the equation (1.3.2) entails that the elliptic operator  $\Gamma$  annihilates  $T_k^*(g) - \mathcal{T}_k(g) + B_k(g)$ . Since both  $T_k^*(g)$  and  $B_k(g)$  belong to  $\mathcal{H}$  this implies by the general result in § xx that  $\mathcal{T}_k(g)$  belongs to  $\mathcal{H}$  which proves the requested equality (1.4) and at the same time the operator equation

$$(1.4.2) \quad T_k^* = -\mathcal{T}_k(g) + B_k$$

**1.5 An inequality.** Let  $f \in C^\infty(\square)$  and  $\lambda$  is a positive real number. Then

$$\|\mathcal{T}_k(f) - \frac{1}{2}B_k(f) - \lambda \cdot f\|^2 =$$

$$\|\mathcal{T}_k(f) - \frac{1}{2}B_k(f)\|^2 + \lambda^2 \cdot \|f\|^2 - \lambda(\langle \mathcal{T}_k(f) - \frac{1}{2}B_k(f), f \rangle + \langle f, \mathcal{T}_k(f) - \frac{1}{2}B_k(f) \rangle)$$

The last term is  $\lambda$  times

$$(i) \quad \langle \mathcal{T}_k(f), f \rangle + \langle f, \mathcal{T}_k(f) \rangle - \langle f, B_k f \rangle$$

where we used that  $B_k$  is symmetric. Now  $T_k = \mathcal{T}_k$  holds on  $C^\infty(\square)$  and the definition of adjoint operators give

$$(ii) \quad \langle \mathcal{T}_k(f), f \rangle = \langle f, T_k^* \rangle$$

Then (1.4.2) implies that (i) is zero and hence we have proved

$$(iii) \quad \|\mathcal{T}_k(f) - \frac{1}{2}B_k(f) - \lambda \cdot f\|^2 = \lambda^2 \cdot \|f\|^2 + \|\mathcal{T}_k(f) - \frac{1}{2}B_k(f)\|^2 \geq \lambda^2 \cdot \|f\|^2$$

From (iii) and the triangle inequality for norms we obtain

$$(iv) \quad \|T_k(f) - \lambda \cdot f\| \geq \lambda \cdot \|f\| - \frac{1}{2}\|B_k(f)\|$$

Now  $B_k$  has a finite operator norm and if  $\lambda \geq \|B_k\|$  we see that

$$(v) \quad \|T_k(f) - \lambda \cdot f\| \geq \frac{\lambda}{2} \cdot \|f\|$$

Finally, since  $C^\infty(\square)$  is dense in  $\mathcal{D}(T_k)$  it is clear that (v) gives

$$(1.5.1) \quad \|T_k(f) - \lambda \cdot f\| \geq \frac{\lambda}{2} \cdot \|f\| \quad : f \in \mathcal{D}(T_k)$$

### § 2. Proof of Theorem 3.1

Suppose we have found some  $\lambda^* \geq \frac{1}{2} \cdot \|B\|$  such that  $T_k - \lambda$  has a dense range in  $\mathcal{H}$  for every  $\lambda \geq \lambda^*$ . If this is so we fix  $\lambda \geq \lambda^*$  and take some  $g \in \mathcal{H}$ . The hypothesis gives a sequence  $\{f_n \in \mathcal{D}(T_k)\}$  such that

$$\lim_{n \rightarrow \infty} \|T(f_n) - \lambda \cdot f_n - g\| = 0$$

In particular  $\{\|T_k(f_n) - \lambda \cdot f_n\|\}$  is a Cauchy sequence in  $\mathcal{H}$  and (1.5.x) implies that  $\{f_n\}$  is a Cauchy sequence in the Hilbert space  $\mathcal{H}$  and hence converges to a limit  $f_*$ . Since the operator  $T_k$  is closed we conclude that  $f_* \in \mathcal{D}(T)$  and we get the equality

$$T_k(f_*) - \lambda \cdot f_* = g$$

Finally, since the graph of  $T$  is contained in  $T_1$  we have the requested equation

$$P(f_*) - \lambda f_* = g$$

Thus finishes the proof of Theorem 1.6 provided we have established the existence of  $\lambda_*$  above

**2.1 Density of the range.** By the construction of adjoint operators the range of  $T_k - \lambda \cdot E$  fails to be dense if and only if  $T_k^* - \lambda$  has a non-zero kernel. So assume that

$$(i) \quad T_k^*(f) - \lambda \cdot f = 0$$

for some  $f \in \mathcal{D}(T_k^*)$  which is not identically zero. Notice that  $T_k$  sends real-valued functions into real-valued functions. So above we can assume that  $f$  is real-valued and also assume that  $f$  is normalised so that

$$\int_{\square} f^2(x, s) dx ds = 1$$

By (\*\*) the equation (xx) gives

$$(ii) \quad \mathcal{T}_k(f) + \lambda \cdot f - B(f) = 0$$

Let us then consider the function

$$V(s) = \int_0^\pi f^2(x, s) dx$$

Recall from § xx that the  $\mathcal{H}$ -function  $f$  is of class  $C^1$ . Now

$$(iii) \quad \frac{1}{2} \cdot V'(s) = \int_0^\pi f \cdot \frac{\partial f}{\partial s} dx$$

By (ii) we have

$$\frac{\partial f}{\partial s} - a(x) \frac{\partial f}{\partial x} - b \cdot f = B(f) - \lambda \cdot f$$

Hence the right hand side in (iii) becomes

$$-\lambda \cdot V(s) + \int_0^\pi f(x, s) \cdot B(f)(x, s) dx + \int_0^\pi a(x, s) \cdot f(x, s) \cdot \frac{\partial f}{\partial x}(x, s) dx$$

By partial integration the last term is equal to

$$-\frac{1}{2} \int_0^\pi \partial_x(a)(x, s) \cdot f^2(x, s) dx$$

Set

$$M = \frac{1}{2} \cdot \max_{(x,s) \in \square} |\partial_x(a)(x, s)|$$

Then we get the inequality

$$\frac{1}{2} \cdot V'(s) \leq (M - \lambda) \cdot V(s) + \int_0^\pi f(x, s) \cdot B(f)(x, s) dx$$

Set

$$\Phi(s) = \int_0^\pi |f(x, s)| \cdot |B(f)(x, s)| dx$$

Since the  $L^2$ -norm of  $f$  is one the Cauchy-Schwarz inequality gives

$$\int_{-\pi}^\pi \Phi(s) ds \leq \sqrt{\int_\square |B(f)(x, s)|^2 dx ds} \leq \|B(f)\|$$

where the last equality follows since the squared integral of  $B(f)$  is majorized by its squared norm in  $\mathcal{H}$ . When  $\lambda > M$  it follows from (xx) that

$$(\lambda - M) \cdot V(s) + \frac{1}{2} \cdot V'(s) \leq \Phi(s)$$

Next, since  $f$  is double periodic we have  $V(-\pi) = V(\pi)$  so after an integration (xx) gives

$$(\lambda - M) \cdot \int_\pi^\pi V(s) ds = \int_{-\pi}^\pi \Phi(s) ds \leq \|B(f)\|$$

By (xx) we have  $\int_\pi^\pi V(s) ds = 1$  which gives a contradiction if  $\lambda > M + \|B(f)\|$ .

**Remark.** Set

$$\tau = \min_f \|B(f)\|$$

with the minimum taken over functions  $f \in \mathcal{D}(T_0^*)$  whose  $L^2$ -integral is normalised by (xx). The proof has shown that the kernel of  $T_0^* - \lambda$  is zero for all  $\lambda > M + \tau$ .

**A special solution.**

Let  $f(x)$  be a periodic  $C^\infty$ -function on  $[0, \pi]$ . Put

$$Q = a(x, s) \cdot \frac{\partial}{\partial x} + b(x, s)$$

Let  $\eta(s)$  be a  $C^\infty$ -function of  $s$  and  $m$  a positive integer. If  $\lambda > 0$  is a real number, we set

$$(i) \quad g_\lambda(x, s) = \eta(s) \cdot f + \eta(s) \cdot \sum_{j=1}^{j=m} \frac{(s-\pi)^j}{j!} \cdot (Q-\lambda)^j(f) \quad : 0 \leq s \leq \pi$$

We choose  $\eta$  to be a real-valued  $C^\infty$ -function such that  $\eta(s) = 0$  when  $s \leq 1/4$  and  $-1$  if  $s \geq 1/2$ . Hence  $g_\lambda(x, s) = 0$  in (i) when  $0 \leq s \leq 1/4$  and we extend the function to  $[-\pi \leq s \leq \pi]$  where  $g_\lambda(x, -s) = g_\lambda(x, s)$  if  $0 \leq s \leq \pi$ . So now  $g_\lambda$  is  $\pi$ -periodic with respect to  $s$  and vanishes when  $|s| \leq 1/4$ .

**Exercise.** If  $1/2 \leq s \leq \pi$  we have  $\eta(s) = 1$ . Use (i) to show that

$$(P + \lambda)(g_\lambda) = \frac{\partial g_\lambda}{\partial s} - (Q - \lambda)(g_\lambda) = \frac{(s - \pi)^m}{m!} \cdot (Q - \lambda)^{m+1}(f)$$

hold when  $1/2 \leq s \leq \pi$ . At the same time  $g_\lambda(s) = 0$  when  $0 \leq s \leq 1/4$ . So  $(P + \lambda)(g)$  is a function whose derivatives with respect to  $s$  vanish up to order  $m$  at  $s = 0$  and  $s = \pi$  and is therefore doubly periodic of class  $C^m$  in  $\square$ . Now Theorem 2.2 applies. For a given  $k \geq 2$  we choose a sufficiently large  $m$  and find  $h(x, s)$  so that

$$P(h) + \lambda \cdot h = (P + \lambda)(g_\lambda)(x, s)$$

where  $h$  is  $s$ -periodic, i.e.

$$h(x, 0) = h(x, \pi)$$

Notice also that  $g_\lambda(x, 0) = 0$  while  $g_\lambda(x, \pi) = f(x)$ . Set

$$g_*(x) = h - g_\lambda$$

Then  $P(g_*) + \lambda \cdot g_* = 0$  and

$$g_*(x, 0) - g_*(x, \pi) = f(x)$$

Above we started with the  $C^\infty$ -function. Given  $k \geq 2$  we can take  $m$  sufficiently large during the constructions above so that  $g_*$  belongs to  $\mathcal{H}^{(k)}(\square)$ .